國立臺灣大學理學院數學學系

碩士論文

Department of Mathematics College of Science National Taiwan University Master Thesis

麥克斯威爾電磁系統的強唯一連續延拓性及其定量分析

Quantitative uniqueness estimate of strong unique continuation property for the Maxwell system with anisotropic media

孫裕沛

Yu-Pei Sun

指導教授:王振男 博士

Advisor: Jenn-Nan Wang, Ph.D.

中華民國 101 年 6 月

June 2012

誌謝

感謝王振男教授兩年以來的用心指導,以及同儕們平時的幫忙,彼此花時間討論出一些有趣的數學論述以及想 法,讓我可以完成這篇文章,也同時感謝我的家人及麻糬的陪伴跟鼓勵,以後我會更努力的!!!



中文摘要

在這篇文章中,我們考慮非零解在時諧性麥克斯威爾系統的局部行為,其系統為非等向 性的媒體。而我們主要得到的結果是此系統的強連續延拓性在某些條件之下將會成立,並且 導出強連續延拓性的定量分析,也可以得到非零解趨近到零的速度。

我們主要運用到的工具為 Carleman 估計導出 Three-balls 不等式,再運用另一個 Carleman 估計以及 Three-balls 不等式推導出 Doubling 不等式,因此可得出強連續延拓性的定量分析。



中文關鍵字:卡勒門估計,麥克斯威爾系統,非等向性,強唯一連續延拓性

Abstract

In this article, we consider the local behavior of a non-trivial solution for the time-harmonic Maxwell system with anisotropic media. The main result of this article is the bound on the vanishing order of the solution of the Maxwell system, which is a quantitative estimate of the strong unique continuation property(SUCP). And the most important tool is *Carleman* estimate. Our strategy in the proof is to derive doubling inequality through three-balls inequality.



Key words: Carleman estimate, Maxwell system, anisotropic media, strong unique continuation

誌譐	甘	i	Ĺ
中文	て摘要	ii	i
英文	て摘要	iii	i
		介紹1	
第二	章	主要定理	ŀ
第三	章	定理一證明	;
		定理二以及定理三證明10)
參考	亡文獻		,



Quantitative uniqueness estimate of strong unique continuation property for the Maxwell system with anisotropic media

1. Introduction

The Maxwell system is firstly mentioned by James Clerk Maxwell in the paper "On Physical Lines of Force" which is published in 1861. He derived it from Gauss's law, Faraday's law and Ampre's circuital law. Furthermore, he derived electromagnetic wave equations in 1865 and claim that light is an electromagnetic wave. In fact, he established the fundamental electrodynamics and had a significant impact on modern physics.

1. *Gauss's law*: The total electric flux coming out of a closed surface is equal to the total charge enclosed by that closed surface. It means that

2. *Gauss law* for the Magnetic Fields: The total magnetic flux coming out of a closed surface is always zero. It means that

$$\oint \mu_0 H \cdot dA = 0$$

3. *Faraday's law*: The line integral of electric field over a closed contour is equal to the time rate of change of the total magnetic flux that goes through any arbitrary surface that is bounded by the closed contour. It means that

$$\oint E \cdot ds = \frac{\partial}{\partial t} \iint \mu_0 H \cdot dA$$

4. Ampre's circuital law: The line integral of magnetic field over a closed contour is equal to the total current plus the time rate of change of the total electric

flux that goes through any arbitrary surface that is bounded by the closed contour. It means that

$$\oint H \cdot ds = \iint J \cdot dA + \frac{\partial}{\partial t} \iint \epsilon_0 E \cdot dA$$

where E is the electric field, H is the magnetic field, J is the current density, Q is the total charge, ϵ_0 is the permittivity of vacuum, and μ_0 is the permeability of vacuum.

Actually, Maxwell system can describe more complicated physical phenomenon in real life, so it has a general form, which is depending on the medium. Now we assume J = 0 to simplify the problem.

So we can define that $E = (E_1, E_2, E_3)$ is the electric field, $H = (H_1, H_2, H_3)$ is the magnetic field and ω is the frequency in a domain Ω . Denote the time-harmonic Maxwell system with anisotropic media

$$\begin{cases} curl E = -i\omega\mu H \\ curl H = i\omega\varepsilon E \end{cases}$$
 in Ω (1.1)

where Ω is an open subset of \mathbb{R}^3 containing $0, \omega \in \mathbb{C} \setminus \{0\}$, and $\varepsilon(x), \mu(x)$ are two real symmetric matrix-valued and positive-definite functions in Ω satisfying the following property :

(a) $\varepsilon(0) = h\mu(0)$ where h is a constant. (b) $\varepsilon, \mu \in \mathbf{C^2}(\Omega)$

We can reduce the Maxwell system to a weakly coupled second order elliptic system. Denote that

$$\gamma_{jl}^{k} = \begin{cases} 1 , \text{if } (\mathbf{k}, \mathbf{j}, \mathbf{l}) \text{ is an even permutation of } (1, 2, 3) \\ -1, \text{if } (\mathbf{k}, \mathbf{j}, \mathbf{l}) \text{ is an odd permutation of } (1, 2, 3) \\ 0 , \text{otherwise} \end{cases}$$

From the Maxwell system we can obtain that

$$\begin{cases} \partial_k E = \nabla E_k - i\omega\gamma^k\mu H\\ \partial_k H = \nabla H_k + i\omega\gamma^k\varepsilon E\end{cases}$$

By simple calculation and (1.1), we know that

$$\begin{cases} \operatorname{div}(curlH) = 0 \\ \operatorname{div}(curlE) = 0 \end{cases} \Rightarrow \begin{cases} \operatorname{div}(\varepsilon E) = 0 \\ \operatorname{div}(\mu H) = 0 \end{cases}$$

So we have that for k = 1, 2, 3, the following formulas is called (1.2)

$$\begin{cases} 0 = \partial_k div(\varepsilon E) = div(\varepsilon \nabla E_k) + div(\partial_k \varepsilon \cdot E - i\omega \varepsilon \gamma^k \mu H) \\ 0 = div(\mu \nabla H_k) + div(\partial_k \mu \cdot H + i\omega \mu \gamma^k \varepsilon E) \end{cases}$$

Now let $P(x, D) = \sum_{j,k} a_{jk}(x) D_j D_k$ be an elliptic operator in Ω such that $a_{jk}(0)$ is a symmetric and positive-definite matrix and $a_{jk}(x) \in \mathbf{C}^2(\Omega)$, so we can rewrite (1.2)

$$\int_{-\infty}^{\infty} P_1(x,D)E + 2\nabla E \cdot div\varepsilon + E \cdot \tilde{\varepsilon} - \sum_{\substack{k=1\\3}}^{3} div(i\omega\varepsilon\gamma^k\mu H) = 0$$
$$P_2(x,D)H + 2\nabla H \cdot div\mu + E \cdot \tilde{\mu} + \sum_{\substack{k=1\\k=1}}^{3} div(i\omega\mu\gamma^k\varepsilon E) = 0$$

where

$$P_1(x,D) = \sum_{i,j=1}^{3} \varepsilon_{ij}(x) D_i D_j , P_2(x,D) = \sum_{k,l=1}^{3} \mu_{kl}(x) D_k D_l,$$
$$\tilde{\varepsilon} = \sum_{m,n=1}^{3} D_m D_n \varepsilon(x) , \tilde{\mu} = \sum_{m,n=1}^{3} D_m D_n \mu(x)$$

So it implies that

$$\begin{cases} |P_1(x,D)E| \le \alpha_1 |E| + \alpha_2 |\nabla E| + \alpha_3 |\nabla H| \le \alpha_4 |U| + \alpha_5 |\nabla U| \\ |P_2(x,D)H| \le \beta_1 |H| + \beta_2 |\nabla H| + \beta_3 |\nabla E| \le \beta_4 |U| + \beta_5 |\nabla U| \end{cases}$$
(1.3)

where U = (E, H) is the non-trivial solution for the (1.1), α_i, β_i are constants for i = 1, 2, 3, 4, 5, and by (1.3) we can assume that $M_1 = max\{\alpha_4, \alpha_5\}$ and $M_2 = max\{\beta_4, \beta_5\} \Rightarrow$

$$\begin{cases} |P_1(x,D)E| \le M_1|U| + M_2|\nabla U| \\ |P_2(x,D)H| \le M_1|U| + M_2|\nabla U| \end{cases}$$
(1.4)

2. The main theorems

Theorem 1 There exists a positive number $R_1 < 1$ such that if $0 < r_1 < r_2 < r_3 \le R_0$ and $r_1/r_3 < r_2/r_3 < R_1$ then

$$\int_{|x| < r_2} |U|^2 dx \le C \left(\int_{|x| < r_1} |U|^2 dx \right)^{\tau} \left(\int_{|x| < r_3} |U|^2 dx \right)^{1-\tau}$$

for $U = (E, H) \in (L^2(B_{R_0}))^6$ where $B_{R_0} \subset \Omega$ and U is the non-trivial solution for the (1.1), where C depend on r_1/r_3 , r_2/r_3 , $P_1(x, D)$ and $P_2(x, D)$ and $0 < \tau < 1$ is only depending on r_1/r_3 , r_2/r_3 .

And then we want to show the quantitative estimate of strong unique continuation property for the Maxwell system. The strong unique continuation means that

For all
$$U = (E, H) \in H^1_{loc}(\Omega)$$
 vanishes of infinite order at 0, then U=0 in Ω

Theorem 2 gives the upper bound on the vanishing order of the solution of the Maxwell system, and theorem 3 is the quantitative estimate of strong unique continuation property.

Theorem 2 If $U = (E, H) \in (L^2_{loc}(\Omega))^6$ is a non-trivial solution of Maxwell system,

then we can find a constant R_2 depending on $P_1(x, D)$, $P_2(x, D)$ and constant m_1 depending on $P_1(x, D)$, $P_2(x, D)$ and $||U||_{L^2(|x| < R_2^2)}/||U||_{L^2(|x| < R_2^4)}$ satisfying

$$\int_{|x| < R} |U|^2 dx \ge K R^{m_1}$$

where R is sufficient small and the constant K depending on R_2 , U.

Theorem 3 Let $U = (E, H) \in (L^2_{loc}(\Omega))^6$ be a non-trivial solution to the Maxwell system. Then there exists positive constant R_3 and C_3 depending on $P_1(x, D)$, $P_2(x, D)$ and m_1 such that if $0 < r \le R_3$,

$$\int_{|x|<2r} |U|^2 dx \le C_3 \int_{|x|$$

where m_1 is the constant obtained in theorem 2.

3. Proofs

Proof of the theorem 1

First we denote that $\varphi_{\beta} = \varphi_{\beta}(|x|) = exp((\frac{\beta}{2})(\log|x|)^2)$ and recall a Carleman estimate [1]:

Lemma For any $\beta > 0$ large enough. Let S be a small neighborhood of 0, and $u: S \setminus \{0\} \subset \Omega \to \mathbb{R}$ and that $u \in H^2(S \setminus \{0\})$ with compact support. Then we have

$$\beta^{3} \int \varphi_{\beta}^{2} |x|^{-n} |u|^{2} dx + \beta \int \varphi_{\beta}^{2} |x|^{-n+2} |\nabla u|^{2} dx \leq \tilde{C}_{0} \int \varphi_{\beta}^{2} |x|^{-n+4} |P(x,D)u|^{2} dx$$
(1.5)

for some positive constant \tilde{C}_0 depending only on P(x, D). Now ε, μ are \mathbb{C}^2 functions, and $U = (E, H) \in (L^2_{loc}(\Omega))^6$ then $U = (E, H) \in (H^1_{loc}(\Omega))^6$ [2] and using regularization, *Friedrich's Lemma* and ellipticity of P(x, D). We can see that $U = (E, H) \in (H^2_{loc}(\Omega \setminus \{0\}))^6$.

Consider that $0 < r_1 < r_2 < R < 1$, $B_R \subset \Omega$ where R is a constant. Define a cut-off function $\phi(x) \in C_0^{\infty}(\mathbb{R}^n)$ satisfying $0 \le \phi(x) \le 1$ and

$$\phi(x) = \begin{cases} 0, \text{ if } |x| \leq \frac{r_1}{e} \\ 1, \text{ if } \frac{r_1}{2} \leq |x| \leq er_2 \\ 0, \text{ if } |x| \geq 3r_2 \end{cases}$$

where exp(1) = e. And it is easy to know that for all multiindex α and C_1, C_2 are constants.

$$\begin{cases} |D^{\alpha}\phi(x)| \le C_1 r_1^{-|\alpha|}, \quad \forall \quad \frac{r_1}{e} \le |x| \le \frac{r_1}{2} \\ |D^{\alpha}\phi(x)| \le C_2 r_2^{-|\alpha|}, \quad \forall \quad er_2 \le |x| \le 3r_2 \end{cases}$$
(1.6)

We assume n = 3 in the lemma because of the domain $\Omega \in \mathbb{R}^3$ and then apply (1.5) to ϕE and ϕH . Firstly, we consider ϕE and use (1.4),(1.5),(1.6) and Cauchy-Schwarz inequality. We obtain

$$\begin{split} \beta^{3} \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-3} |E|^{2} dx + \beta \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-1} |\nabla E|^{2} dx \\ & \leq \beta^{3} \int \varphi_{\beta}^{2} |x|^{-3} |\phi E|^{2} dx + \beta \int \varphi_{\beta}^{2} |x|^{-1} |\nabla \phi E|^{2} dx \end{split}$$

$$\leq \tilde{C}_{0} \int \varphi_{\beta}^{2} |x| |P_{1}(x, D)(\phi E)|^{2} dx$$

$$\leq \tilde{C}_{0} \left\{ \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x| (2M_{1}^{2} |\phi U|^{2} + 2M_{2}^{2} |\phi \nabla U|^{2}) dx$$

$$+ \int_{r_{1}/e < |x| < r_{1}/2} \varphi_{\beta}^{2} |x|^{-3} (C_{1} |U|^{2} + C_{2} |x|^{2} |\nabla U|^{2}) dx$$

$$+ \int_{er_{2} < |x| < 3r_{2}} \varphi_{\beta}^{2} |x|^{-3} (C_{1} |U|^{2} + C_{2} |x|^{2} |\nabla U|^{2}) dx \right\}$$

$$\leq \tilde{C}_{1} \left\{ \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-3} (|U|^{2} + |x|^{-1} |\nabla U|^{2}) dx$$

$$+ \int_{r_{1}/e < |x| < r_{1}/2} \varphi_{\beta}^{2} |x|^{-3} (|U|^{2} + |x|^{2} |\nabla U|^{2}) dx$$

$$+ \int_{er_{2} < |x| < 3r_{2}} \varphi_{\beta}^{2} |x|^{-3} (|U|^{2} + |x|^{2} |\nabla U|^{2}) dx$$

$$+ r_{1}^{-3} \varphi_{\beta}^{2} (r_{1}/e) \int_{r_{1}/e < |x| < r_{1}/2} (|U|^{2} + |x|^{2} |\nabla U|^{2}) dx$$

$$+ r_{2}^{-3} \varphi_{\beta}^{2} (er_{2}) \int_{er_{2} < |x| < 3r_{2}} (|U|^{2} + |x|^{2} |\nabla U|^{2}) dx$$

$$+ r_{2}^{-3} \varphi_{\beta}^{2} (er_{2}) \int_{er_{2} < |x| < 3r_{2}} (|U|^{2} + |x|^{2} |\nabla U|^{2}) dx$$

$$(1.7)$$

where $\tilde{C}_1 = max\{2\tilde{C}_0M_1^2, 2\tilde{C}_0M_2^2, \tilde{C}_0C_1, \tilde{C}_0C_2\}$ and $\tilde{C}_1e^3 = \tilde{C}_2$ We introduce a corollary in [3]

Corollary For $0 < a_3 < a_1 < a_2 < a_4$ such that $B_{a_4r} \subset \Omega$, we can show the following inequality

$$\int_{a_1r < |x| < a_2r} ||x|^{|\alpha|} D^{\alpha} u|^2 dx \le C' \int_{a_3r < |x| < a_4r} |u|^2 dx$$

where C' is a constant independent of r and $|\alpha| \leq 2$. So by the corollary, it implies (1.8)

$$\begin{split} \beta^{3} \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-3} |E|^{2} dx + \beta \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-1} |\nabla E|^{2} dx \\ & \leq \tilde{C}_{3} \bigg\{ \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} (|x|^{-3} |U|^{2} + |x|^{-1} |\nabla U|^{2}) dx \\ & + r_{1}^{-3} \varphi_{\beta}^{2} (r_{1}/e) \int_{r_{1}/4 < |x| < r_{1}} |U|^{2} dx + r_{2}^{-3} \varphi_{\beta}^{2} (er_{2}) \int_{2r_{2} < |x| < 4r_{2}} |U|^{2} dx \bigg\} \end{split}$$

where $\tilde{C}_1, \tilde{C}_2, \tilde{C}_3$ are independent of r_1, r_2 .

And we have the same conclusion for the ϕH , so we obtained

$$\beta^3 \int_{r_1/2 < |x| < er_2} \varphi_\beta^2 |x|^{-3} |H|^2 dx + \beta \int_{r_1/2 < |x| < er_2} \varphi_\beta^2 |x|^{-1} |\nabla H|^2 dx$$

$$\leq \tilde{C}_{3} \left\{ \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2}(|x|^{-3}|U|^{2} + |x|^{-1}|\nabla U|^{2}) dx + r_{1}^{-3} \varphi_{\beta}^{2}(r_{1}/e) \int_{r_{1}/4 < |x| < r_{1}} |U|^{2} dx + r_{2}^{-3} \varphi_{\beta}^{2}(er_{2}) \int_{2r_{2} < |x| < 4r_{2}} |U|^{2} dx \right\}$$

$$(1.9)$$

Therefore, we can combine the inequality (1.8) and (1.9) such that

$$\begin{split} \beta^{3} \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-3} (|E|^{2} + |H|^{2}) dx + \beta \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-1} (|\nabla E|^{2} + |\nabla H|^{2}) dx \\ &= \beta^{3} \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-3} |U|^{2} dx + \beta \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} |x|^{-1} |\nabla U|^{2} dx \\ &\leq 2 \tilde{C}_{3} \bigg\{ \int_{r_{1}/2 < |x| < er_{2}} \varphi_{\beta}^{2} (|x|^{-3} |U|^{2} + |x|^{-1} |\nabla U|^{2}) dx \\ &+ r_{1}^{-3} \varphi_{\beta}^{2} (r_{1}/e) \int_{r_{1}/4 < |x| < r_{1}} |U|^{2} dx + r_{2}^{-3} \varphi_{\beta}^{2} (er_{2}) \int_{2r_{2} < |x| < 4r_{2}} |U|^{2} dx \bigg\} \end{split}$$

Now let $\beta_0 \ge 1$ and $\beta^3 \ge \beta \ge \beta_0 \ge 3\tilde{C}_3$, then we can get another inequality (1.10)

$$\begin{split} \int_{r_1/2 < |x| < er_2} \varphi_{\beta}^2 |x|^{-3} |U|^2 dx + \int_{r_1/2 < |x| < er_2} \varphi_{\beta}^2 |x|^{-1} |\nabla U|^2 dx \\ & \leq \tilde{C}_4 \left\{ r_1^{-3} \varphi_{\beta}^2(r_1/e) \int_{r_1/4 < |x| < r_1} |U|^2 dx + r_2^{-3} \varphi_{\beta}^2(er_2) \int_{2r_2 < |x| < 4r_2} |U|^2 dx \right\} \\ & \text{where } \tilde{C} = \frac{1}{\tilde{C}} \left\{ \tilde{C}_4 \left\{ r_1^{-3} \varphi_{\beta}^2(r_1/e) \int_{r_1/4 < |x| < r_1} |U|^2 dx + r_2^{-3} \varphi_{\beta}^2(er_2) \int_{2r_2 < |x| < 4r_2} |U|^2 dx \right\} \end{split}$$

where $\tilde{C}_4 = 1/\tilde{C}_3$, and it is easy to get that

$$\begin{aligned} r_2^{-3}\varphi_{\beta}^2(r_2) \int_{r_1/2 < |x| < r_2} |U|^2 dx &\leq \int_{r_1/2 < |x| < r_2} \varphi_{\beta}^2 |x|^{-3} |U|^2 dx \leq \int_{r_1/2 < |x| < er_2} \varphi_{\beta}^2 |x|^{-3} |U|^2 dx \\ &\leq \tilde{C}_4 \left\{ r_1^{-3} \varphi_{\beta}^2(r_1/e) \int_{r_1/4 < |x| < r_1} |U|^2 dx + r_2^{-3} \varphi_{\beta}^2(er_2) \int_{2r_2 < |x| < 4r_2} |U|^2 dx \right\} \end{aligned}$$

Dividing the term $r_2^{-3}\varphi_{\beta}^2(r_2)$, we obtain

$$\int_{r_{1}/2 < |x| < r_{2}} |U|^{2} dx \leq \tilde{C}_{4} \left\{ (r_{2}/r_{1})^{3} [\varphi_{\beta}^{2}(r_{1}/e)/\varphi_{\beta}^{2}(r_{2})] \int_{r_{1}/4 < |x| < r_{1}} |U|^{2} dx + [\varphi_{\beta}^{2}(er_{2})/\varphi_{\beta}^{2}(r_{2})] \int_{2r_{2} < |x| < 4r_{2}} |U|^{2} dx \right\} \\
\leq \tilde{C}_{5} \left\{ (r_{2}/r_{1})^{3} [\varphi_{\beta}^{2}(r_{1}/e)/\varphi_{\beta}^{2}(r_{2})] \int_{|x| < r_{1}} |U|^{2} dx + (r_{2}/r_{1})^{3} [\varphi_{\beta}^{2}(er_{2})/\varphi_{\beta}^{2}(r_{2})] \int_{|x| < 4r_{2}} |U|^{2} dx \right\}$$
(1.11)

where $\tilde{C}_5 = max\{\tilde{C}_4, 1\}$

By choosing such \tilde{C}_5 , we know that

$$\tilde{C}_5(r_2/r_1)^3[\varphi_\beta^2(r_1/e)/\varphi_\beta^2(r_2)] > 1$$

for $0 < r_1 < r_2 \le 1$.

$$\begin{aligned} &\operatorname{Adding} \int_{|x| < r_1/2} |U|^2 dx \text{ to the both sides of (1.11) and } r_2 < 1/4, \text{ and then we have} \\ &\int_{|x| < r_2} |U|^2 dx \leq \tilde{C}_5 \left\{ (r_2/r_1)^3 [\varphi_{\beta}^2(r_1/e)/\varphi_{\beta}^2(r_2)] \int_{|x| < 1} |U|^2 dx \\ &\quad + (r_2/r_1)^3 [\varphi_{\beta}^2(er_2)/\varphi_{\beta}^2(r_2)] \int_{|x| < r_1/2} |U|^2 dx \right\} \\ &\quad + \tilde{C}_5 (r_2/r_1)^3 [\varphi_{\beta}^2(r_1/e)/\varphi_{\beta}^2(r_2)] \int_{|x| < r_1/2} |U|^2 dx \\ &\leq 2\tilde{C}_5 \left\{ (r_2/r_1)^3 [\varphi_{\beta}^2(r_1/e)/\varphi_{\beta}^2(r_2)] \int_{|x| < r_1} |U|^2 dx \\ &\quad + (r_2/r_1)^3 [\varphi_{\beta}^2(er_2)/\varphi_{\beta}^2(r_2)] \int_{|x| < 1} |U|^2 dx \right\} \end{aligned}$$

Assume $A = (\log r_1 - 1)^2 - (\log r_2)^2$, $B = -1 - 2\log r_2$, and A > 0, B > 0 by simple computation. Therefore, the above inequality becomes

$$\int_{|x| < r_2} |U|^2 dx \le 2\tilde{C}_5 (r_2/r_1)^3 \left\{ \exp(A\beta) \int_{|x| < r_1} |U|^2 dx + \exp(-\beta B) \int_{|x| < 1} |U|^2 dx \right\}$$
(1.12)

By standard argument, we consider two cases Case1: If $exp(A\beta_0) \int_{|x| < r_1} |U|^2 dx < exp(-\beta_0 B) \int_{|x| < 1} |U|^2 dx$ and pick $\beta > \beta_0$ such that that

$$exp(A\beta) \int_{|x| < r_1} |U|^2 dx = exp(-\beta B) \int_{|x| < 1} |U|^2 dx$$

so we have the following important inequality

$$\begin{split} \int_{|x|$$

By the arguments, we can take $C_6 = max \{ exp(\beta_0 B), 4C_5(r_2/r_1)^3 \}$ and get that

$$\int_{|x| < r_2} |U|^2 dx \le \tilde{C}_6 \left(\int_{|x| < r_1} |U|^2 dx \right)^{\frac{B}{A+B}} \left(\int_{|x| < 1} |U|^2 dx \right)^{\frac{A}{A+B}}$$
(1.13)

For the general case, we can assume that $R_1 \leq 1/4$ and $0 < r_1 < r_2 < r_3 \leq R_0$ with $r_1/r_3 < r_2/r_3 \leq 1/4$.

By scaling, $\tilde{U}(y) = U(r_3 y)$, $\tilde{\varepsilon_{ij}}(y) = \varepsilon_{ij}(r_3 y)$, $\tilde{\mu_{ij}}(y) = \mu_{ij}(r_3 y)$ We can have the same conclusion by above argument and obtain

$$\int_{|y| < r_2/r_3} |\tilde{U}|^2 dx \le C \bigg(\int_{|y| < r_1/r_3} |\tilde{U}|^2 dx \bigg)^{\tau} \bigg(\int_{|y| < 1} |\tilde{U}|^2 dx \bigg)^{1-\tau}$$
(1.14)

where $\tau = B/(A+B)$ and $C = max \{ exp(\beta_0 B), 4\tilde{C}_5(r_2/r_1)^3 \}$

$$\begin{cases} A = \left(\log(r_1/r_3) - 1\right)^2 - \left(\log(r_2/r_3)\right)^2 \\ B = -1 - 2\log(r_2/r_3) \end{cases}$$

Providing $r_3 < 1$ and \tilde{C}_5 can be chosen independent of r_3 . So undoing the change of variable of (1.14), we have

$$\int_{|x| < r_2} |U|^2 dx \le C \left(\int_{|x| < r_1} |U|^2 dx \right)^{\tau} \left(\int_{|x| < r_3} |U|^2 dx \right)^{1-\tau}$$
(1.15)

100

The proof is now complete.

And then we are going to prove that the Maxwell system have the strong unique continuation property, so we have to prove the two theorems by using theorem 1.

Proof of the theorem 2 and theorem 3

Without loss of generality, we can use the change of coordinates and property (a) to obtain that

$$P_1(0,D) = \sum_{i,j=1}^{3} \varepsilon(0) D_i D_j = \Delta$$
$$P_2(0,D) = \sum_{i,j=1}^{3} \mu(0) D_i D_j = \frac{P_1(0,D)}{h} = \frac{\Delta}{h}$$

So we recall another Carleman estimate[1]: For any $u \in H^2_{loc}(\mathbb{R}^n \setminus \{0\})$ with compact support and for any $m \in \{j + \frac{1}{2} | j \in \mathbb{N}\}$ we have that

$$\sum_{|\alpha| \le 2} \int m^{2-2|\alpha|} |x|^{-2m+2|\alpha|-n} |D^{\alpha}u|^2 dx \le C \int |x|^{-2m+4-n} |\Delta u|^2 dx \tag{2.1}$$

where C only depends on the dimension n.

And from the previous description, we know that $U = (E, H) \in (H^2_{loc}(\Omega \setminus \{0\}))^6$, so we can use the Carleman estimate for U.

Define a cut-off function $\chi(x) \in C_0^\infty(\mathbb{R}^n)$ satisfying $0 \le \chi(x) \le 1$ and

$$\chi(x) = \begin{cases} 0, \text{ if } |x| \le \frac{\delta}{3} \\ 1, \text{ if } \frac{\delta}{2} \le |x| \le \frac{(R_0 + 1)R_0R}{4} = r_4R \\ 0, \text{ if } |x| \ge 2r_4R \end{cases}$$

where $\delta \leq R_0^2 R/4$, $R_0 > 0$ is a small number and it will be determined later, and R is sufficiently small satisfying $0 < R \leq R_0$. Using the (2.1) for χE and χH . Now for χE , we can derive that

$$\sum_{|\alpha| \le 2} m^{2-2|\alpha|} \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}E|^2 dx$$

$$\le \sum_{|\alpha| \le 2} m^{2-2|\alpha|} \int |x|^{-2m+2|\alpha|-3} |D^{\alpha}(\chi E)|^2 dx$$

$$\le C \int |x|^{-2m+1} |\Delta(\chi E)|^2 dx$$

$$\le C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\Delta E|^2 dx + C \int_{|x| > r_4 R} |x|^{-2m+1} |\Delta(\chi E)|^2 dx$$

$$+ C \int_{\delta/3 \le |x| \le \delta/2} |x|^{-2m+1} |\Delta(\chi E)|^2 dx \qquad (2.2)$$

On the right hand side of (2.2), the first term we use the triangle inequality (2.3)

$$C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\Delta E - P_1(x, D)E + P_1(x, D)E|^2 dx$$

$$\le C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\Delta E - P_1(x, D)E|^2 dx$$

$$+ C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |P_1(x, D)E|^2 dx$$

and the first term of the right hand side of (2.3), we can find out that

$$C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\Delta E - P_1(x, D)E|^2 dx$$

$$= C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |(P_1(0, D) - P_1(x, D))E|^2 dx$$

$$= C \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\sum_{i,j=1}^3 (\varepsilon_{ij}(0) - \varepsilon_{ij}(x))D_i D_j E|^2 dx$$

$$\le C \int_{\delta/2 \le |x| \le r_4 R} (|x| sup |\varepsilon'_{ij}(x)|)^2 |x|^{-2m+1} \sum_{i,j=1}^3 |D_i D_j E|^2 dx$$

$$\le C' \sum_{\alpha=2} r_4^2 R^2 \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |D^\alpha E|^2 dx$$
(2.4)

since $\varepsilon_{ij}(x)$ is **C²**-function and C, C' are constants.

So by (2.2),(2.3),(2.4) and (1.4) we obtain

$$\sum_{|\alpha|\leq 2} m^{2-2|\alpha|} \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}E|^2 dx$$

$$\leq C' \sum_{\alpha=2} r_4^2 R^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |D^{\alpha}E|^2 dx$$

$$+2CM_1^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |U|^2 dx + 2CM_2^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |\nabla U|^2 dx$$

$$+C \int_{|x|>r_4 R} |x|^{-2m+1} |\Delta(\chi E)|^2 dx + C \int_{\delta/3\leq |x|\leq \delta/2} |x|^{-2m+1} |\Delta(\chi E)|^2 dx$$
(2.5)
And we can have the same argument for χH to get that

And we can have the same argument for χH to get that

$$\begin{split} \sum_{|\alpha|\leq 2} m^{2-2|\alpha|} \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}H|^2 dx \\ &\leq \tilde{C}' \sum_{\alpha=2} r_4^2 R^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |D^{\alpha}H|^2 dx \\ &+ 2\tilde{C}M_1^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |U|^2 dx + 2\tilde{C}M_2^2 \int_{\delta/2\leq |x|\leq r_4 R} |x|^{-2m+1} |\nabla U|^2 dx \\ &+ \tilde{C} \int_{|x|> r_4 R} |x|^{-2m+1} |\Delta(\chi H)|^2 dx + \tilde{C} \int_{\delta/3\leq |x|\leq \delta/2} |x|^{-2m+1} |\Delta(\chi H)|^2 dx (2.6) \end{split}$$
And then we can derive (2.7) from (2.5) (2.6)

And then we can derive (2.7) from (2.5),(2.6)

$$\begin{split} \sum_{|\alpha| \le 2} m^{2-2|\alpha|} \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}U|^2 dx \\ & \le C'' \sum_{\alpha=2} r_4^2 R^2 \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |D^{\alpha}U|^2 dx \\ & + C_M \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |U|^2 dx + C_M \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+1} |\nabla U|^2 dx \\ & + \hat{C} \int_{|x| > r_4 R} |x|^{-2m+1} |\Delta(\chi U)|^2 dx + \hat{C} \int_{\delta/3 \le |x| \le \delta/2} |x|^{-2m+1} |\Delta(\chi U)|^2 dx \ (2.7) \end{split}$$
where $C'' = C' + \tilde{C}', \ C_M = 2((C + \tilde{C})M_1^2 + (C + \tilde{C})M_2^2)$ and $\hat{C} = C + \tilde{C}$

By choosing

$$R = \frac{1}{m\sqrt{C''}} \qquad \text{and} \qquad r_4^2 R^2 = \frac{R_0^2 (R_0 + 1)^2}{16m^2 C''}$$

Choosing $R_0 < 1$ (such that $\frac{R_0^2(R_0+1)^2}{16} < \frac{1}{2}$) and $m = m(R_0)$ large enough such that

$$\sum_{|\alpha| \le 2} m^{2-2|\alpha|} \int_{\delta/2 \le |x| \le r_4 R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}U|^2 dx$$
$$\le 2C \int_{\delta/3 \le |x| \le \delta/2} |x|^{-2m+1} |\Delta(\chi U)|^2 dx + 2C \int_{r_4 R < |x| < 2r_4 R} |x|^{-2m+1} |\Delta(\chi U)|^2 dx (2.8)$$

The first three terms on right hand side of (2.7) is absorbed by the left hand side when the |x| is small enough.

And then by the definition of χ , it is easy to obtain that for all multiindex α

$$\begin{cases} |D^{\alpha}\chi(x)| \leq C_{3}\delta^{-|\alpha|}, \quad \forall \quad \frac{\delta}{3} \leq |x| \leq \frac{\delta}{2} \\ |D^{\alpha}\chi(x)| \leq C_{4}(r_{4}R)^{-|\alpha|}, \quad \forall \quad r_{4}R \leq |x| \leq 2r_{4}R \end{cases}$$
(2.9)

where C_3, C_4 are constants. Now we provide $R_0 \leq 1/16$ such that $R_0^2 \leq r_4$, so by the corollary in theorem 1 and (2.9), we can derive (2.10) from (2.8)

$$m^{2}(2\delta)^{-2m-3} \int_{\delta/2 \leq |x| \leq 2\delta} |U|^{2} dx + m^{2} (R_{0}^{2}R)^{-2m-3} \int_{2\delta < |x| \leq R_{0}^{2}R} |U|^{2} dx$$

$$\leq \sum_{|\alpha| \leq 2} m^{2-2|\alpha|} \int_{\delta/2 \leq |x| \leq r_{4}R} |x|^{-2m+2|\alpha|-3} |D^{\alpha}U|^{2} dx$$

$$\leq 2CC_{3} \sum_{|\alpha| \leq 2} \delta^{-4+2|\alpha|} \int_{\delta/3 \leq |x| \leq \delta/2} |x|^{-2m+1} |D^{\alpha}U|^{2} dx$$

$$+ 2CC_{4} \sum_{|\alpha| \leq 2} (r_{4}R)^{-4+2|\alpha|} \int_{r_{4}R < |x| < 2r_{4}R} |x|^{-2m+1} |D^{\alpha}U|^{2} dx$$

$$\leq C_{3}'(\delta/3)^{-2m-3} \int_{|x| \leq \delta} |U|^{2} dx + C_{4}'(r_{4}R)^{-2m-3} \int_{|x| \leq R_{0}R} |U|^{2} dx \quad (2.10)$$

where C'_3, C'_4 are independent of R_0, R and m.

Adding $m^2(2\delta)^{2m-3} \int_{|x|<\delta/2} |U|^2 dx$ to both sides of (2.10), we can get the left hand side

$$\begin{split} m^{2}(2\delta)^{-2m-3} &\int_{|x| \le 2\delta} |U|^{2} dx + m^{2} (R_{0}^{2}R)^{-2m-3} \int_{2\delta < |x| \le R_{0}^{2}R} |U|^{2} dx \\ &= \frac{1}{2} m^{2} (2\delta)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx + \frac{1}{2} m^{2} (2\delta)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx \\ &+ m^{2} (R_{0}^{2}R)^{-2m-3} \int_{2\delta < |x| \le R_{0}^{2}R} |U|^{2} dx \\ &\geq \frac{1}{2} m^{2} (2\delta)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx + m^{2} (R_{0}^{2}R)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx \\ &+ m^{2} (R_{0}^{2}R)^{-2m-3} \int_{2\delta < |x| \le R_{0}^{2}R} |U|^{2} dx \\ &= \frac{1}{2} m^{2} (2\delta)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx + m^{2} (R_{0}^{2}R)^{-2m-3} \int_{|x| \le 2\delta} |U|^{2} dx (2.11) \end{split}$$

Combine (2.10) and (2.11), it implies that

$$\frac{1}{2}m^{2}(2\delta)^{-2m-3}\int_{|x|\leq 2\delta}|U|^{2}dx + m^{2}(R_{0}^{2}R)^{-2m-3}\int_{|x|\leq R_{0}^{2}R}|U|^{2}dx$$

$$\leq (C_{3}'+m^{2})(\delta/3)^{-2m-3}\int_{|x|\leq \delta}|U|^{2}dx + C_{4}'(r_{4}R)^{-2m-3}\int_{|x|\leq R_{0}R}|U|^{2}dx$$
We can rewrite $C_{4}'(r,R)^{-2m-3}$ and set $m^{2}(R^{2}R)^{-2m-3}C_{4}'m^{-2}(R_{0}^{2})^{2m+3}$

Teller.

1.2.2

We can rewrite $C'_4(r_4R)$ (r_4^{-}) Therefore,

$$C_4'm^{-2}(\frac{R_0^2}{r_4})^{2m+3} = C_4'm^{-2}(\frac{4R_0}{R_0+1})^{2m+3} \le C_4'm^{-2}(4R_0)^{2m+3} \le exp(-2m)$$

For all
$$R_0 \leq 1/16$$
 and $m^2 \geq C'_4$. Thus, we can derive

$$\frac{1}{2}m^2(2\delta)^{-2m-3} \int_{|x|\leq 2\delta} |U|^2 dx + m^2(R_0^2 R)^{-2m-3} \int_{|x|\leq R_0^2 R} |U|^2 dx$$

$$\leq (C'_3 + m^2)(\delta/3)^{-2m-3} \int_{|x|\leq \delta} |U|^2 dx + m^2(R_0^2 R)^{-2m-3} e^{-2m} \int_{|x|\leq R_0 R} |U|^2 dx$$
The above argument is valid for all $m = m(R_0) = i + 1/2$ with $i \in \mathbb{N}$ and m large

The above argument is valid for all $m = m(R_0) = j + 1/2$ with $j \in \mathbb{N}$, and m large enough. So we can assume a j_0 which depends on R_0 , and $R_j = (\sqrt{C''}(j+1/2))^{-1}$. For all $j \ge j_0$, we have the following inequality (2.12)

$$\frac{1}{2}m^{2}(2\delta)^{-2m-3}\int_{|x|\leq 2\delta}|U|^{2}dx + m^{2}(R_{0}^{2}R_{j})^{-2m-3}\int_{|x|\leq R_{0}^{2}R_{j}}|U|^{2}dx \\
\leq (C_{3}'+m^{2})(\delta/3)^{-2m-3}\int_{|x|\leq \delta}|U|^{2}dx + m^{2}(R_{0}^{2}R_{j})^{-2m-3}e^{\frac{-2c}{R_{j}}}\int_{|x|\leq R_{0}R_{j}}|U|^{2}dx \\
\text{where } c = \frac{1}{\sqrt{cm}}$$

 $\sqrt{C'}$

We can observe that $R_{j+1} < R_j < 2R_{j+1}$ for all $j \in \mathbb{N}$ by simple calculation. If we can find a R such that $R_{j+1} < R \leq R_j$, it implies $R_0 R_j \leq 2R_{j+1}/16 \leq R_{j+1}$ since $R_0 < 1/16$ and then we have such relation $R_0 R_j \le R_{j+1} < R < R_j < 2R_{j+1}$. So we can get a conclusion

$$\begin{cases}
\int_{|x| \le R_0^2 R} |U|^2 dx \le \int_{|x| \le R_0^2 R_j} |U|^2 dx \\
e^{\frac{-2c}{R_j}} \int_{|x| \le R_0 R_j} |U|^2 dx \le e^{\frac{-c}{R}} \int_{|x| \le R} |U|^2 dx
\end{cases}$$
(2.13)

If there exists a $s \in \mathbb{N}$ such that

$$R_{j+1} < R_0^{2s} \le R_j \text{ for some } j \ge j_0 \tag{2.14}$$

It is just replacing R by R_0^{2s} on the above description. By (2.12) and (2.13), we can obtain (2.15)

$$\frac{1}{2}m^{2}(2\delta)^{-2m-3}\int_{|x|\leq 2\delta}|U|^{2}dx + m^{2}(R_{0}^{2}R_{j})^{-2m-3}\int_{|x|\leq R_{0}^{2s+2}}|U|^{2}dx \\
\leq (C_{3}'+m^{2})(\delta/3)^{-2m-3}\int_{|x|\leq \delta}|U|^{2}dx + m^{2}(R_{0}^{2}R_{j})^{-2m-3}exp(-cR_{0}^{-2s})\int_{|x|\leq R_{0}^{2s}}|U|^{2}dx$$

So now our goal is to find a appropriate s and R_0 to claim an inequality which is

$$exp(-cR_0^{-2s})\int_{|x| \le R_0^{2s}} |U|^2 dx \le \frac{1}{2} \int_{|x| \le R_0^{2s+2}} |U|^2 dx$$
(2.16)

By theorem 1, we assume $r_1 = R_0^{2s+2}, r_2 = R_0^{2s}$ and $r_3 = R_0^{2s-2}$ where $s \ge 1$. And $r_1/r_3 < r_2/r_3 \le R_0^2 \le 1/4$, then we divide $\left(\int_{|x| < R_0^{2s}} |U|^2 dx\right)^{1-\tau}$ to both sides of (1.15)

$$\left(\int_{|x|
and it implies (2.17)$$

$$\int_{|x|< R_0^{2s}} |U|^2 dx \Big/ \int_{|x|< R_0^{2s+2}} |U|^2 dx \le \overline{C}^{1/\tau} \bigg(\int_{|x|< R_0^{2s-2}} |U|^2 dx \Big/ \int_{|x|< R_0^{2s}} |U|^2 dx \bigg)^a$$

where $\overline{C} = max \{ exp(\beta_0(-1 - 4logR_0)), 4\tilde{C}_5(R_0)^{-6} \}$

 $a = \frac{1-\tau}{\tau} = \frac{(4\log R_0 - 1)^2 - (2\log R_0)^2}{-1 - 4\log R_0}$ by definition of τ in the proof of theorem 1.

We can see that

$$\begin{cases}
1 < \overline{C} \le \tilde{C}_5 R_0^{-\beta_1} \\
2 < a \le -5 log R_0
\end{cases}$$
(2.18)

where $\beta_1 = max\{6, 4\beta_0\}.$

Because $exp(\beta_0(-1 - 4\log R_0)) = e^{-\beta_0}R_0^{-4\beta_0}$ and $\tilde{C}_5 \ge 1$, we have first inequality. For the second inequality, we consider that

$$a = \frac{-4\log R_0 + 2 - 1/4\log R_0 + \log R_0}{1/4\log R_0 + 1} > -3\log R_0 + 2 > 2$$
$$a < \frac{-3\log R_0 - \log R_0}{1 - 1/5} < -5\log R_0$$

And we use (2.17) recursively

$$\int_{|x|$$

For all $s \in \mathbb{N}$. And by definition of a we know that $\tau = 1/(a+1)$, from (2.18) we have

$$\frac{a^{s-1}-1}{\tau(a-1)} = \frac{(a+1)(a^{s-1}-1)}{a-1} \le 3a^{s-1}$$

Then we derive the following inequality from (2.19)

$$\begin{split} \int_{|x|$$

Multiply
$$exp(-cR_0^{-2s})$$
 on both sides, we obtain (2.20)
 $exp(-cR_0^{-2s}) \int_{|x|< R_0^{2s}} |U|^2 dx \Big/ \int_{|x|< R_0^{2s+2}} |U|^2 dx$
 $\leq exp(-cR_0^{-2s}) (\tilde{C}_5^3 R_0^{-3\beta_1})^{(-5\log R_0)^{s-1}} \left(\int_{|x|< R_0^2} |U|^2 dx \Big/ \int_{|x|< R_0^4} |U|^2 dx \right)^{a^{s-1}}$

Let $\kappa = -\log R_0$, and compute $\log(\tilde{C}_5^3 R_0^{-3\beta_1})^{(5\kappa)^{s-1}} = (5\kappa)^{s-1} (\log \tilde{C}_5^3 + 3\beta_1 \kappa)$. So we can find out that if R_0 small enough, it means that κ sufficient large, then we

have (2.21)

$$\frac{cR_0^{-2s}}{4} > (5\kappa)^{s-1} (\log \tilde{C}_5^3 + 3\beta_1 \kappa) \Rightarrow (\tilde{C}_5^3 R_0^{-3\beta_1})^{(5\kappa)^{s-1}} < exp(\frac{cR_0^{-2s}}{4}) < \frac{1}{2} exp(\frac{cR_0^{-2s}}{2})$$
The (2.21) holds for all $s \in \mathbb{N}$, and now we should fix R_0 such that (2.21) holds and

the $m(R_0)$ and $j(R_0)$ are fixed as well. Fixing R_0 , we can derive from (2.20)

$$exp(-cR_0^{-2s}) \int_{|x|< R_0^{2s}} |U|^2 dx \le \frac{1}{2} exp(\frac{-cR_0^{-2s}}{2}) \left(\int_{|x|< R_0^2} |U|^2 dx \middle/ \int_{|x|< R_0^4} |U|^2 dx \right)^{a^{s-1}} \times \int_{|x|< R_0^{2s+2}} |U|^2 dx$$

Our goal is (2.16), so coping with the term $\left(\int_{|x|< R_0^2} |U|^2 dx / \int_{|x|< R_0^4} |U|^2 dx\right)^d$ If we have an estimate that

$$\left(\int_{|x|$$

Then (2.16) is proved for appropriate s. So now we have to find s such that (2.22) holds.

Assume
$$N = \left(\int_{|x| < R_0^2} |U|^2 dx \middle/ \int_{|x| < R_0^4} |U|^2 dx \right)$$
, take loglog for (2.22), we have
 $log2 - log(ac) + loglogN \le s(-2logR_0 - loga)$

By (2.18), we know $-2logR_0 - loga > 0$ for all $R_0 \le 1/16$, so we can define a number s_0 as

$$s_0 = \min\{s \in \mathbb{N} | s \ge (\log 2 - \log(ac) + \log \log N)(-2\log R_0 - \log a)^{-1}\}$$

So the claim (2.16) holds for all $s \ge s_0$.

But now s should also be chosen to assure (2.14) holds.

Let s_1 be the smallest positive integer such that $R_0^{2s_1} \leq R_{j_0}$, then we can find a $j_1 \in \mathbb{N}$ with $j_1 \geq j_0$ such that $R_{j_1+1} < R_0^{2s_1} \leq R_{j_1}$. We now can define s_p depending on $P_1(x, D)$, $P_2(x, D)$ and N as

$$s_p = max\{s_0, s_1\}$$

For this s_p , (2.14),(2.21),(2.22) hold. Thus, we set $m = j_1 + \frac{1}{2}$ and $m_1 = 3 + 2m$ plus into (2.15).

$$\left(\frac{m_{1}-3}{8}\right)^{2}(2\delta)^{-2\left(\frac{m_{1}-3}{2}\right)-3} \int_{|x|\leq 2\delta} |U|^{2} dx + \left(\frac{m_{1}-3}{8}\right)^{2} (R_{0}^{2}R_{j_{1}})^{-2\left(\frac{m_{1}-3}{2}\right)-3} \int_{|x|\leq R_{0}^{2s_{p}+2}} |U|^{2} dx \\
\leq (C_{3}' + \left(\frac{m_{1}-3}{2}\right)^{2})(\delta/3)^{-2\left(\frac{m_{1}-3}{2}\right)-3} \int_{|x|\leq \delta} |U|^{2} dx \quad (2.23)$$

So consider the second term of the left hand side of (2.23)

$$\frac{1}{8}(m_1-3)^2 (R_0^2 R_{j_1})^{-m_1} \int_{|x| \le R_0^{2s_p+2}} |U|^2 dx \le (C_3' + (\frac{m_1-3}{2})^2)(\delta/3)^{-m_1} \int_{|x| \le \delta} |U|^2 dx$$

then it implies

$$\frac{(m_1-3)^2}{8C'_3+2(m_1-3)^2} (3R_0^2 R_{j_1})^{-m_1} \int_{|x| \le R_0^{2s_p+2}} |U|^2 dx \le (\delta)^{-m_1} \int_{|x| \le \delta} |U|^2 dx$$
(2.24)

(2.24) is valid for all $\delta \leq R_0^{2s_p+2}/4$ because of the definition of δ . So the proof of theorem 2 is complete with $R_2 = R_0$. And the first term of the left hand side of (2.23)

$$\frac{1}{8}(m_1 - 3)^2 (2\delta)^{-m_1} \int_{|x| \le 2\delta} |U|^2 dx \le (C'_3 + (\frac{m_1 - 3}{2})^2)(\delta/3)^{-m_1} \int_{|x| \le \delta} |U|^2 dx$$

then it implies

$$\int_{|x| \le 2\delta} |U|^2 dx \le \frac{8C_3' + 2(m_1 - 3)^2}{(m_1 - 3)^2} 6^{m_1} \int_{|x| \le \delta} |U|^2 dx \tag{2.25}$$

(2.25) is valid for all $\delta \leq R_0^{2s_p+2}/4$ because of the definition of δ . So the proof of theorem 3 is complete with $R_3 = R_0^{2s_p+2}/4$ and $C_3 = \frac{8C'_3 + 2(m_1 - 3)^2}{(m_1 - 3)^2} 6^{m_1}$.



References

- REGBAOUI, R.: Strong uniqueness for second order differential operators. J.Differential Equations 141(1997), no. 2,201-217.
- [2] C. Weber, Regularity theorems for Maxwell equation, Math. Methods Appl. Sci., 3(1981),523-536.
- [3] HORMANDER, L.: The analysis of linear partial differential operators.III. Grundlehren der Mathematischen Wissenschaften 274. Springer-Verlag, Berlin, 1985.
- [4] CHING-LUNG LIN, GEN NAKAMURA and JENN-NAN WANG: Quantitative uniqueness for second order elliptic operators with strongly singular coefficients, Rev. MAT. IBEROAMERICANA 27(2011), no 2, 475-491
- [5] TU NGUYEN and JENN-NAN WANG: Quantitative uniqueness estimate for the Maxwell system with Lipschitz anisotropic media, Proc. Amer. Math. Soc. 140 (2012), 595-605